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Variations in Snow Crystal Riming and Z_{DR}:

A Case Analysis

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Abstract

A case study in terms of variations in differential reflectivity (Z_{DR}) observed at X-band and snow crystal riming is presented for a light snow event that occurred near Greeley, Colorado on 26–27 November 2015. In the early portion of the event, Z_{DR} values at near surface levels were low (~ 0 to 0.25 dB). During a second time period approximately eight hours later, Z_{DR} values became distinctly positive ($\sim +2$ to +3 dB). Digital photographs of the snow particles were obtained by a Multi-Angle Snowflake Camera (MASC) installed at a range of 13 km from the radar. Image processing and machine learning techniques applied to the MASC data showed that the snow particles were more heavily rimed during the low Z_{DR} time period. The aerodynamic effects of these rime deposits promoted a wider distribution of hydrometeor canting angles. The shift toward more random particle orientations underlies the observed reduction in Z_{DR} during the period where more heavily rimed particles were observed in the MASC data.

1. Introduction

Synoptic scale mid-latitude precipitation systems often have vertical extents that include much of the troposphere. This depth spans a wide temperature range; the organized upward air motions in these systems also promote significant supercooled liquid water concentrations. These conditions support several hydrometeor development processes including active dendritic growth near the -15°C level (Kennedy and Rutledge 2011), and areas of local riming enhancement and secondary ice production (Giangrande et al. 2016). These cloud microphysical complexities, and especially the variability in the microwave backscattering properties of different ice particle types as they undergo riming, contribute to the challenge of detecting riming conditions using conventionally scanning dual-polarization radars (Vogel et al. 2015).

40 Historically, differential reflectivity (Z_{DR}) has been shown to provide useful information
41 on the drop size distributions as well as shape and orientation characteristics of raindrops (Seliga
42 and Bringi 1976; Bringi et al. 1998). The interpretation of the Z_{DR} observed from snow is more
43 complicated. Unlike liquid drops, frozen hydrometeors do not have well-defined size – shape
44 relationships. Nevertheless, a large number of polarimetric radar measurements show fascinating
45 signatures especially related to positive Z_{DR} and K_{DP} aloft (e.g., Matrosov et al. 2001; Kennedy
46 and Rutledge 2011; Kumjian et al. 2013; Ryzhkov et al. 2016; Williams et al. 2015). Frozen
47 precipitation particles generally exhibit larger spatial orientation angle fluctuations than rain
48 drops (Ryzhkov et al. 2011; Matrosov et al. 2005), and furthermore the scattering properties of
49 ice crystals can be dependent on the intricate, detailed particle structures (e.g., Botta et al. 2011;
50 Kuo et al. 2016; Lu et al. 2016). The refractive index properties of snowflakes also vary
51 according to the amount of air contained in the overall ice structure (Smith 1984). When riming
52 deposits frozen cloud droplets on the outer surface of snow particles, snowflake shape, fall
53 orientation, and density characteristics are all altered, which in turn complicates the
54 interpretation of polarimetric radar data. It follows that in-situ measurements of these particle
55 characteristics are important but difficult to obtain at the required resolution. Aircraft-based
56 imaging probes yield important information on particle size and shape (crystal habit), as well as
57 qualitative indication of riming and environmental conditions favorable for crystal formation, but
58 the imaging is generally in one plane only and as a result complete orientation information is
59 typically not available (Wolde and Vali 2001). Under horizontally homogeneous conditions, the
60 elevation angle dependence of polarimetric radar data can reveal information of ice crystal
61 orientation or “flutter” (Matrosov et al. 2005).

62 This study is motivated by the physical reasons underpinning observed differential
63 reflectivity (Z_{DR}) variations on the order of 2 dB that occurred near the ground by a 9 GHz (X-
64 band) radar during a multi-hour, widespread, light snow event in the high plains of Colorado.
65 During this event, digital photographs of the snow particles were obtained by a Multi-Angle
66 Snowflake Camera (MASC; Garrett et al. 2012) that was located at a range of 13 km from the
67 radar. Image processing techniques that automatically extract a variety of solid hydrometeor
68 characteristics from MASC data sets have recently been developed (Praz et al. 2017). In addition
69 to the classification of particles by habit, these image processing procedures also provide
70 information on the axis ratios, orientation angles, and riming level of the imaged particles.
71 Garrett et al. (2015) use a simpler approach to determine the statistics of aspect ratio and
72 orientation angle of snow, rimed snow, and graupel based on MASC data. Using herein the Praz
73 et al. (2017) methodology, MASC-based snow particle characteristics are examined from two
74 time periods that displayed appreciably different Z_{DR} regimes. During the low Z_{DR} time period,
75 the MASC image analyses showed that consistently higher levels of riming existed. These rimed
76 particles, especially at larger diameters, also had a broader distribution of orientation angles and
77 slightly larger (more spherical) axis ratios. These snow particle orientation characteristics are
78 consistent with the observed lower Z_{DR} levels. During the contrasting 2 to 3 dB Z_{DR} time period,
79 riming was less evident in the MASC images and the particles maintained more horizontal
80 orientation angles.

81 This paper is organized as follows. Section 2 describes the instrumentation site and
82 briefly the Multi-Angle Snowflake Camera system and the CSU-CHILL radar's X-band system
83 which is used herein, as well as the larger scale meteorology of the event which occurred on 26-
84 27 November 2015. Section 3 describes in detail the radar and MASC data for the two periods

85 characterized by different Z_{DR} signatures, as well as microphysical characteristics which form the
86 central core of this study. Finally, we end with discussion and conclusions sections.

87

88 **2. Instrumentation and Meteorological Situation**

89 *a) Easton Site Instrumentation*

90 The data for this analysis were collected as a part of the MASC + Radar project
91 (MASCRAID; Notaros et al. 2016). The MASCRAID project involved the installation of a ground
92 instrumentation site at the Easton Valley View Airport (Fig. 1). The South Platte River valley is
93 located between the CSU-CHILL radar and the Easton site. The reduced terrain heights in the
94 river valley combined with higher terrain elevations at Easton vs. CSU-CHILL (1432 m vs. 1460
95 m) allow clutter-free radar observations to be made at antenna elevations as low as 0.9° over
96 Easton. At this elevation angle, the 0.33° wide CSU-CHILL X-band main beam illuminates a
97 region located between 150 and 224 m above ground level over the Easton instrumentation site
98 (13 km from the radar on an azimuth of 171°). To reduce the impact of horizontal wind on the
99 precipitation observations, a 2/3-scale (8 m outer diameter) Double Fence Intercomparison
100 Reference (DFIR) wind screen was constructed at Easton within which MASC and several other
101 ground instruments were installed. The DFIR enabled substantial reduction in the horizontal
102 winds (Notaros et al. 2016).

103 The basic MASC design includes three computer-controlled digital cameras located at
104 36° angular intervals around a horizontal plane. Hydrometeors falling towards the common
105 viewing area of these three cameras pass through two infrared motion detection beams. The
106 interruption of these beams triggers the cameras and their associated flash illumination systems.
107 In the Colorado State University (CSU) MASC, the three horizontal plane cameras are Unibrain

108 980 digital camera. Each of these cameras generates 5 Mpixel files with a maximum resolution
109 of 35 μm . The measurement area of the optical system is 18 cm^2 . To improve 3D reconstructions
110 of the imaged particles, two additional downward looking, lower-resolution cameras were added
111 (Fig. 1 inset). While the processing of images collected by all five cameras can be done, only
112 images from the three horizontal plane cameras that are the basis of the conventional MASC
113 design were used in this analysis.

114 The MASC image pixels are recorded using 256 grayscale levels. This intensity data
115 permits the extraction of various texture-based descriptors which are then used to automatically
116 detect frozen cloud drops that accumulate on the outer surface of solid hydrometeors during
117 riming, following the procedure introduced by Praz et al. (2017).

118

119 *b) The CSU-CHILL Radar*

120 The radar data presented here was collected with the CSU-CHILL operating in dual
121 frequency mode (Junyent et al. 2015). In this configuration, a two frequency antenna feed allows
122 dual polarization measurements to be made simultaneously at both 3 GHz (S-band) and 9 GHz
123 (X-band). The same antenna reflector system is used at both frequencies, yielding a main (3 dB)
124 beam width of 1.0° at S-band and 0.33° at X-band.

125 All of the radar data in this analysis was obtained from the CSU-CHILL X-band channel.
126 This system uses a 25 KW peak power magnetron transmitter whose output is split, causing the
127 antenna to simultaneously radiate both horizontally and vertically polarized waves. The
128 reflectivity measurements obtained with the CSU-CHILL X-band system averaged within 1 dB
129 of the values that were independently observed by a vertically-pointed, X-band continuous wave
130 radar (Precipitation Occurrence Sensing System; Sheppard 1990) that was operating at Easton

131 during the analyzed event. The CSU-CHILL Z_{DR} calibration was based on vertically pointed
132 scans that were done at selected times when light precipitation was in progress at the radar site
133 during the MASCRAD project. The Z_{DR} accuracy is estimated to be within ± 0.12 dB.

134

135 *c) Meteorological Situation*

136 Salient features of the surface conditions during this event are summarized by the time
137 history of the data collected at the Greeley-Weld County (KGXY) Automated Weather
138 Observing System (AWOS) located ~11 km from the Easton site (Fig. 2). The color-coded sets
139 of vertical lines in each panel of Fig. 2 mark the two time periods when different Z_{DR} regimes
140 were noted in the immediate Easton area: During period 1 (26 November 2015, 18-20 UTC), Z_{DR}
141 values averaged $\sim +0.2$ dB; in period 2 (27 November 2015, 02-0430 UTC), Z_{DR} was
142 consistently more positive, averaging $\sim +2.2$ dB. Reflectivity levels were ~ 10 dBZ in period 1
143 and ~ 4 dBZ in period 2. The precipitation that occurred during this event was related to a
144 shallow, anticyclonic upslope flow that developed the day after an initial surface cold front that
145 passed the Greeley Airport near 1430 UTC on the morning of 25 November 2015. The
146 accumulation rate of snow that occurs in this synoptic environment is generally light (Rasmussen
147 et al. 1992). During period 1, surface winds were from the east-northeast at speeds of $4\text{--}5$ ms^{-1}
148 (Fig. 2a and b). These persistent synoptic scale upslope post-frontal surface winds resulted in the
149 development of low clouds and periodic light snow in the Greeley area by time period 1. The
150 cloud cover and low-level cold air advection kept surface temperatures at -6°C during most of
151 the daylight hours (Fig. 2c). Around 21 UTC, wind directions became more northerly and speeds
152 increased slightly to ~ 6 ms^{-1} (Fig. 2a and b). Two instances of gusting winds were reported in
153 the 22–23 UTC period. The temperature and dew point traces show the arrival of colder, drier air

154 associated with this wind shift (Fig. 2c). By the time of period 2, wind speeds decreased as a
155 high pressure system moved into the area (Fig. 2b and d).

156 The overall evolution of the surface data suggests that during the earlier times, including
157 period 1, the combination of upslope flow and subsequently enhanced convergence associated
158 with the weak cold air surge that arrived around 21–22 UTC promoted the development of
159 supercooled liquid water (Rasmussen et al., 1995). After ~ 00 UTC, upward air motions probably
160 decreased as surface winds became light and colder, drier, high pressure conditions developed.

161

162 **3. Data and Analyses**

163 *a) Radar Data*

164 During the MASCRAD project, a prescribed sequence of scans that were focused on the
165 Easton site were conducted by the CSU-CHILL radar. This sequence included a narrow (50°
166 wide) Plan Position Indicator (PPI) volume with two elevation angles (0.9° and 1.5°). These
167 were the lowest elevation angles that were free of ground clutter in the Easton area at S- and X-
168 bands, respectively. Two sweep Range Height Indicator (RHI) scans were also done on azimuths
169 that flanked the Easton site (171° and 172°; see Fig. 1). These scans were repeated at 3 minute
170 intervals. Figure 3 shows representative X-band reflectivity and differential reflectivity (Z_{DR})
171 data collected during RHI scans in time periods 1 and 2 as identified in Fig. 2. In accordance
172 with the generally weak synoptic scale forcing associated with the anticyclonic upslope
173 environment, the echo depth was less than 2 km at both times (Fig. 3a and c). These echo top
174 heights were also consistent with the base of a temperature inversion in the operational sounding
175 data (not shown) from Denver, located ~70 km south of Easton. Reflectivities increased towards
176 the ground at 1851 UTC during period 1, reaching 10–12 dBZ within 200 m of the surface at

177 Easton. In contrast, during period 2 at 0258 UTC, near-surface reflectivities generally remained
178 below 6 dBZ. The Z_{DR} regimes were also quite different at these two times. At 1851 UTC, low-
179 level Z_{DR} near Easton was very close to 0 dB; at 0258 UTC, the entire echo column within ± 1
180 km of Easton contained Z_{DR} values of +3 dB or more (Fig. 3b and d).

181 The higher near-surface reflectivity and reduced Z_{DR} levels observed during period 1
182 indicate that the ice particles were larger and appearing more “spherical” as they descended to
183 the ground. In contrast, during period 2, distinctly positive Z_{DR} values were observed in a low
184 reflectivity environment. Hydrometeor classification schemes based on X-band radar data have
185 associated these low reflectivity – positive Z_{DR} echo characteristics at subfreezing temperatures
186 with the presence of relatively pristine planar ice crystals (Dolan and Rutledge 2009).

187 A more complete summary of the reflectivity and Z_{DR} contrasts observed between
188 periods 1 and 2 is provided by the scatterplot shown in Figure 4. The data points were extracted
189 from all of the 0.9° PPI sweeps done during these two time periods. Specifically, data from the
190 range gates on this PPI scan surface that were located within a ± 1 km range interval and a $\pm 1^\circ$
191 azimuth interval of Easton are plotted. During period 1, Z_{DR} remained near 0 dB despite Z_h
192 variations of ~ 15 dB. In period 2, Z_{DR} was more positive, with the highest Z_{DR} tending to be
193 associated with lowest reflectivities. Figure 4 also includes example MASC images from these
194 two time periods. More riming is visibly evident on the crystal photographed during the higher
195 reflectivity / low Z_{DR} conditions of period 1.

196

197 *b) MASC Data*

198 Image processing techniques to classify individual MASC snow particle images have
199 recently been developed (Praz et al. 2017). In this effort, a regularized Multinomial Logistic

200 Regression (MLR) model was developed based on a training data set of approximately 3500
201 MASC images that had been manually classified with regard to particle type (among six
202 predefined classes), presence of melting, and degree of riming. The classification was automated
203 based on the distribution of 25 geometrical (e.g., particle maximum dimension, aspect ratio,
204 shape complexity, and fractal dimension) and textural (e.g., mean brightness, local interpixel
205 variability, and co-occurrence matrix) descriptors. For each individual image, the MLR model
206 assigns a probability of belonging to each predefined class as well as a riming index, as
207 illustrated in Fig. 5. These probabilities are then weighted over the three views of the same
208 particle in order to assign a unique label for each hydrometeor.

209 The riming index is an indicator of the extent of riming on the surface of the particle. It is
210 defined on a scale from zero to one with 0 indicating no cloud droplets visible on the particle
211 surface and 1 indicating a graupel particle (i.e., complete obscuration of the initial particle by an
212 extensive rime accumulation). The riming index assignment was based on a training data set in
213 which the visually apparent degree of riming was put into one of five categories following
214 Mosimann et al. (1994).

215 To relate the CHILL X-band Z_{DR} to the MASC-based riming index, all Z_{DR} range gate
216 data within an area over and surrounding the MASCRAD site were extracted from the PPI
217 sweeps, and a standard time-filtering procedure was applied (Lee 1980) in order to reduce the
218 gate-to-gate fluctuations. Such filtering techniques had been shown to be in very good agreement
219 with time-series of FIR (finite impulse response) range-filtered Z_{DR} (for example, see Fig. 2 of
220 Thurai et al. 2012). The filtered Z_{DR} values were then time-interpolated to match the time of the
221 MASC images. Finally, these filtered data were used to generate the riming index versus X-band
222 Z_{DR} scatter plot shown in Fig. 6. Also shown in this plot are the means and standard deviations

223 for the Z_{DR} and riming index data. Riming index values were higher during time period 1. The
224 average riming index during period 1 (0.53) corresponds to continuous coating of cloud droplets
225 on the frozen hydrometeor surfaces, while the 0.30 average value during period 2 indicates that
226 the cloud droplet coverage was 50% or less.

227 The MASC images also provide information on particle aspect ratio (minor axis length /
228 major axis length) and fall orientation. Both of these factors affect Z_{DR} (Bringi and Chandrasekar
229 2001). Orientation is usually described by the zenith angle (θ) and the azimuthal angle (ϕ) of the
230 symmetry axis relative the local vertical Z-axis. For scattering analysis, the canting angle (β),
231 which is the angle between the projection of the symmetry axis onto the polarization plane and
232 the projection of the vertical direction (Z-axis) onto this same plane (Bringi and Chandrasekar
233 2001 Fig. 2.6a), is more relevant. At horizontal incidence, say along the X-axis and restricting
234 the symmetry axis to be in the YZ plane, the canting angle is representative of θ . To make a
235 correspondence with MASC definition of orientation angle, the vertical Z-axis is taken to be
236 orthogonal to the horizontal plane (XY plane) where the cameras are located. For instance, the
237 central camera can be considered to be located along the X-axis so the image from that camera
238 will be in the YZ-plane. The orientation angle of the ellipse in this plane can be considered to be
239 the canting angle. In the method of Praz et al. (2017), aspect ratio is based on the minor and
240 major lengths of the ellipse that yields the least squares best fit to the particle's outer edge.
241 Orientation angle is the rotation angle between major axis of this ellipse and the horizontal plane
242 (between -90° and 90°). Henceforth, we will not distinguish between the orientation angle from
243 the MASC and the canting angle relevant to the radar, and will frequently use “flutter” to
244 characterize snow crystal fall mode. Finally, projected area is defined as the number of pixels
245 included within a single particle's image.

246 Figure 7a shows the standard deviation of the orientation angle as a function of the
247 minimum projected area for time periods 1 and 2, respectively. The X-axis values in this plot
248 indicate the lowest projected area value for which the standard deviation of canting angle was
249 calculated (for example, the standard deviation of canting angle associated with the 20 mm² X-
250 axis value is based on the orientation angle data for all particles with a projected area of 20 mm²
251 or larger). This method was used to separate the orientation angle fluctuations associated with
252 the larger sized particles that provide greater contributions to the X-band radar returns. For
253 projected areas larger than ~16 mm², the standard deviation of canting angle is larger in period 1;
254 this period also had a relatively high riming index. The difference in the standard deviation of
255 canting angle between these two periods approached a factor of 2 (~50° in period 1 vs. ~25° in
256 period 2) for projected areas larger than ~28 mm². The DFIR wind screen, including the 2/3 scale
257 design used at Easton, has been shown to achieve a snow collection efficiency of ~90% under
258 the generally light (< 6 ms⁻¹) surface wind speeds (Rasmussen et al. 2012) that prevailed during
259 this event. Based on this overall effectiveness of the DFIR wind screen design, we believe that
260 these MASC-derived particle orientation angles were not significantly affected by instrument-
261 induced turbulence. The increased standard deviation of orientation angle observed during period
262 1 is in qualitative agreement with the reduced Z_{DR} levels observed during this period (see Fig. 11
263 later in the text).

264 Additional characteristics of the orientation angle distributions are shown in Fig. 7b.
265 This plot shows the mean and standard deviation of the canting angles associated with the same
266 sequence of minimum projected area values that were used in Fig. 7a. For both periods 1 and 2,
267 the magnitude of the mean canting angle generally was small (< 20°). (When averaging was
268 done over all of the particle sizes, the magnitude of the mean canting angle was under 3° during

269 both time periods). The wider distribution of canting angles, especially for projected areas larger
270 than $\sim 25 \text{ mm}^2$, during the riming conditions of time period 1 is evident.

271 The hydrometeor axis ratios derived from the MASC images recorded during analysis
272 time period 1 and 2 are shown in Fig. 8. The axis ratio mean and standard deviation values are
273 calculated and presented following the procedures that were used in Fig. 7b. While the axis ratios
274 were consistently larger (more spherical) during period 1, the axis ratio differences between the
275 two time periods cannot be considered to be statistically significant. In contrast to the standard
276 deviation of canting angle distributions shown in Fig. 7b, significantly wider distributions of axis
277 ratios were found during period 2. The reduced riming during period 2 allowed the more pristine
278 ice crystals to retain their intrinsically low axis ratios. The presence of these low axis ratio
279 particles broadened the ice particle axis ratio distribution.

280 The results of the MASC-based hydrometeor classification scheme were also examined
281 for periods 1 and 2. As depicted in Fig. 5, the hydrometeor identification scheme presented in
282 Praz et al. (2017) resolves six particle classifications. During the analyzed periods on 26–27
283 November 2015, no classifications in the combination of columns and plates category were
284 made. A summary of the classification results, combined with the previously described riming
285 index results, is presented in Fig. 9. The most frequently identified classifications in both time
286 periods were planar crystals and aggregates. Graupel identifications, indicative of the most
287 advanced stages of riming, were only observed during period 1 (Fig. 9a). This is in agreement
288 with the generally higher riming index values that were generally found on all types of particles
289 in period 1 relative to period 2.

290 Figure 9b shows the MASC image complexity value associated with each of the
291 hydrometeor classes for time periods 1 and 2. The complexity number (similar to Garrett et al.

292 2012) is defined as the ratio of the perimeter length of the silhouette of the MASC particle image
293 to the circumference of a circle of the same cross sectional area. Thus, a circular image results in
294 a complexity value of 1.0; images with complicated perimeters, like branched ice crystals, etc.
295 will have complexities greater than three. Complexity was uniformly higher during period 2 (low
296 riming); in contrast, complexity was reduced in period 1 (more active riming). The higher riming
297 indices recorded during period 1 are consistent with a more extensive rime coating that would
298 produce smoother particle edges and fill in the ice structural irregularities (Moisseev et al. 2017).

299

300 **4. Discussion**

301 The event considered here involved a shallower, more simply organized echo system than
302 the deep cyclonic type winter precipitation systems noted in Rasmussen et al (1992). Particle
303 aggregation processes were not dominant in the sense that the MASC hydrometeor
304 classifications consistently contained a significant fraction of individual ice crystals (Fig. 9). At
305 the $\sim -12^{\circ}\text{C}$ echo top temperature level, growth of planar type ice crystals was favored (Bailey
306 and Hallett 2009). The prevalence of such crystals is consistent with the tendency for more
307 positive Z_{DR} values to be observed near echo top level (Williams et al. 2015).

308 During period 1, the slightly stronger upslope (easterly component) flow and pre-cold
309 surge convergence at low levels were more favorable for the development of upward air motions
310 and supercooled cloud droplets necessary for riming. The combination of increased riming index
311 (0.5–0.6), lower complexity (< 2), and large standard deviation of canting angles of the larger
312 particles (around 50°) during period 1 is consistent with rimed crystals and aggregates exhibiting
313 complicated fall modes probably generated by spatial asymmetries in the growing rime deposit
314 (Zikmunda and Vali 1972; Jayaweera and Mason 1965). This causes Z_{DR} to reduce dramatically

315 (from the unrimed case) to near 0 dB. This interpretation differs from conventional arguments
316 that increased riming of ice crystals during the early stages causes the aspect ratio to increase
317 (i.e., more spherical) resulting in decrease of the Z_{DR} (Moisseev et al. 2017).

318 In contrast, during period 2, the riming index is much reduced (0.2–0.3), with large
319 complexity (> 3), and much lower standard deviation of canting angles (20–30°) which are
320 consistent with low flutter of the unrimed crystals leading to the conventionally observed higher
321 Z_{DR} values in the range 2–3 dB (e.g., Matrosov et al. 2005; Williams et al. 2015). Figure 10
322 shows a schematic representation of the differences in the snow particle characteristics between
323 periods 1 and 2 with the proviso that the mean aspect ratios inferred from MASC data are not
324 significantly different between periods 1 and 2.

325 The premise that increased standard deviation of canting angles or flutter for plate-like
326 crystals tends to decrease Z_{DR} is well-founded via scattering calculations using spheroid models
327 (e.g., Bringi and Chandrasekar 2001; Melnikov and Straka 2013). An analytic equation derived
328 by Jameson (1987) and, later, in a simpler manner by Bringi and Chandrasekar (2001) shows the
329 essential features of this dependence which is more complex and involves the copolar correlation
330 coefficient as well as linear depolarization ratio (LDR) (see Appendix). The assumptions are
331 spheroid shapes with mean canting angle ≈ 0 and further for simplicity the copolar correlation
332 coefficient ≈ 1 (the mean canting angle assumption is in agreement with MASC data in Fig. 7a
333 whilst Melnikov and Straka 2013 show that the correlation coefficient is typically > 0.95 for
334 plate-like crystals exhibiting flutter). Fig. 11 depicts the modeled behavior of Z_{DR} versus σ_{β} for
335 two LDR values. The two points marked show $[Z_{DR}, \sigma_{\beta}]$ for the two periods where Z_{DR} is based
336 on averaged radar measurements whilst σ_{β} is based on MASC data in Fig. 7a for projected area $>$
337 30 mm^2 since the canting of the largest particles dominates the radar returns. The points

338 representing periods 1 and 2 are on curves of constant but different LDR values, respectively,
339 -32 and -20 dB. As mentioned in the Appendix, the sensitivity of the S-band cross-polar
340 receiver precluded measurement of such low LDR signals. (The CSU-CHILL X-band system
341 operates in simultaneous H and V transmission mode and thus cannot measure LDR). The higher
342 LDR of -20 dB inferred during period 2 is consistent with [$Z_{DR} = 2.5$ dB; $\sigma_{\beta} = 20^{\circ}$] and Eq.
343 (A1). One reason for the high LDR during period 2 is likely related to the very large standard
344 deviation in aspect ratios as shown in Fig. 8, especially for projected area $> 20\text{--}30$ mm².
345 Equation (7.39) in Bringi and Chandrasekar (2001) shows that for a given σ_{β} and mean aspect
346 ratio, the LDR can be enhanced by the variance of axis ratios which is not included in Eq. (A1).
347 Such large values of LDR due to pristine dendrites exhibiting flutter have been measured, though
348 at Ka-band, by Matrosov et al. (2005), who found values between -20 to -25 dB at low
349 elevation angles. During period 1, the much lower LDR of -32 dB is consistent with [$Z_{DR} =$
350 0.5 dB; $\sigma_{\beta} = 50^{\circ}$] and Eq. (A1). However, the standard deviation of axis ratios is much lower (see
351 Fig. 8) especially for projected area > 30 mm² so that no enhancement of LDR due to this feature
352 is expected.

353

354 **5. Conclusions**

355 The shallow-echo precipitation episode that affected the MASCRAD project area on 26–
356 27 November 2015 provided a light snow regime where the aggregation process was restrained
357 enough to preserve a detectable population of planar and columnar ice crystals based on the
358 MASC particle classifications. Over the course of the event, the MASC image data showed
359 significantly different degrees of particle riming between two time periods that contained
360 contrasting Z_{DR} regimes. The MASC-based particle orientation and axis ratio measurements

361 revealed the manner in which the riming affected hydrometeor physical properties that impact
362 Z_{DR} . Specifically, the standard deviation of the hydrometeor canting angles was observed to be
363 significantly larger during the time period when riming was more evident. These more random
364 particle orientations contributed to the observed 1.5–2 dB reduction in Z_{DR} .

365 The relationship between Z_{DR} and hydrometeor riming is complicated (Moisseev et al.
366 2017; Vogel et al. 2015). Results differing from those obtained in this work should be expected
367 in other meteorological and microphysical situations. It is anticipated that the high resolution
368 particle photographs collected by the MASC, especially in conjunction with automated image
369 processing techniques, along with other optical instruments that measure fall speeds and particle
370 size distributions at high resolutions, will become increasingly useful. The measurements of
371 these frozen hydrometeor physical characteristics are relevant to explain or model the interesting
372 polarimetric radar signatures observed under different environmental conditions.

373

374

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378 input that significantly improved the paper.

379

380

APPENDIX

381

Z_{DR} and Standard Deviation of Canting Angle

382 To examine the relationship between Z_{DR} and standard deviation of canting angle, we
383 make use of the equations provided in section 7.1.3 in Bringi and Chandrasekar (2001). In

384 particular, equation (7.40a) can be used to derive an approximate variation between the two
 385 parameters. The equation can be rewritten as follows:

$$386 \quad L = \frac{1}{4}(1 - \rho_4) \left\{ 1 - \frac{1}{\sqrt{z_{\text{dr}}}} \right\}^2 \quad (\text{A1})$$

387 where

$$388 \quad \rho_4 = e^{-8\sigma_\beta^2}$$

389 and L is the linear depolarization ratio (LDR), z_{dr} is Z_{DR} in linear units, and σ_β is the standard
 390 deviation of canting angles (the mean canting angle is assumed to be ≈ 0). Note that equation
 391 (A1) is applicable only when the copolar correlation coefficient approaches unity, and
 392 furthermore under the assumption of spheroidal shapes.

393 Equation (A1) can be rewritten as:

$$394 \quad \frac{4L}{\left\{ 1 - \frac{1}{\sqrt{z_{\text{dr}}}} \right\}^2} = 1 - \rho_4$$

395 which in turn can be recast as:

$$396 \quad z_{\text{dr}} = \left\{ \frac{1}{1 - \sqrt{\frac{4L}{1 - \rho_4}}} \right\}^2 \quad (\text{A2})$$

397 In other words, the relationship between Z_{DR} and σ_β is dependent on L.

398 If we now input some example values for Z_{DR} (from CHILL X-band data) and for σ_β (determined
 399 from the MASC images) for the time periods, we obtain the following:

- 400 - For period 1, $Z_{\text{DR}} = 0.5$ dB; $\sigma_\beta = 60^\circ$; calculated LDR ≈ -30 dB
- 401 - For period 2, $Z_{\text{DR}} = 2.5$ dB; $\sigma_\beta = 20^\circ$; calculated LDR ≈ -20 dB

402 These are the two $[Z_{\text{DR}}, \sigma_\beta]$ point pairs that are superimposed onto the two curves in Fig. 11.

403 Note that the cross-polar S-band receiver sensitivity at the range to the Easton measurement site

404 is not sufficient to measure LDR values as low as -20 dB under these weak copolar echo
405 conditions.

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552

553 **Figure Caption List**

554

555 **Figure 1:** (a) MASCRAD instrumentation network and terrain height contours (m MSL) in the
556 immediate CSU-CHILL radar – Easton area. The radar azimuths that flank the Easton site are
557 shown in grey. (b) Overview of the ground instrument installation at Easton. (c) Close-up view
558 of the MASC instrument as modified by CSU. The automated surface weather observations
559 plotted in Fig. 2 were collected at the location marked KGXY AWOS, ~11 km from the MASC
560 instrument at Easton.

561

562 **Figure 2:** Time history of Greeley Airport METAR data. Analysis time periods 1 and 2 are
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564

565 **Figure 3:** Example RHI data: Period 1 (a and b); period 2 (c and d). Red “E” marked along the
566 range axis is the Easton site.

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568 **Figure 4:** Extracted CHILL X-band Z and Z_{DR} data. Sample MASC images are shown for each
569 of the two periods.

570

571 **Figure 5:** MASC automated data processing overview. The procedure is detailed in Section 3b.

572 While the CSU MASC contains 5 cameras, the input to the automated processing system was

573 restricted to images obtained from the three horizontal plane cameras.

574

575 **Figure 6:** Scatterplot of riming index derived from MASC particle image data vs. CSU-CHILL

576 X-band Z_{DR} together with their corresponding means and standard deviations. The Z_{DR} data was

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581 **Figure 7:** (a) Standard deviation of orientation angle vs. minimum projected area threshold for

582 the MASC data collected during period 1 (purple) and period 2 (orange). Values based on less

583 than 10 particle images are plotted in gray. (b) Means (dots) and standard deviations (bars) for

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586 **Figure 8:** Means (dots) and standard deviations (bars) for the hydrometeor axis ratios observed
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588 Fig. 7b. Data from time period 1 is plotted in purple; time period 2 is plotted in orange.

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594 color coded by time period, along the abscissa indicate the number of individual MASC particle
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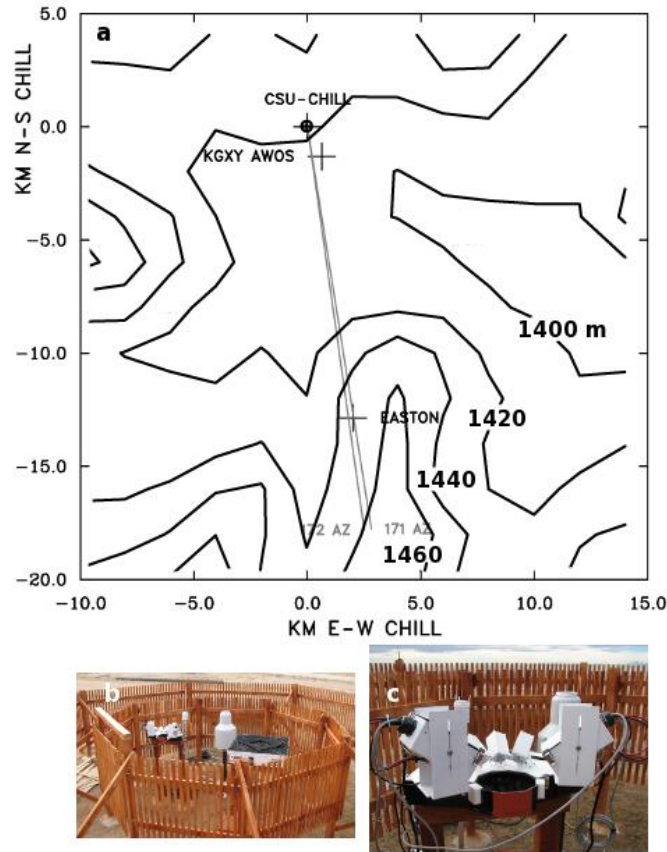
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603

604

Figures

605



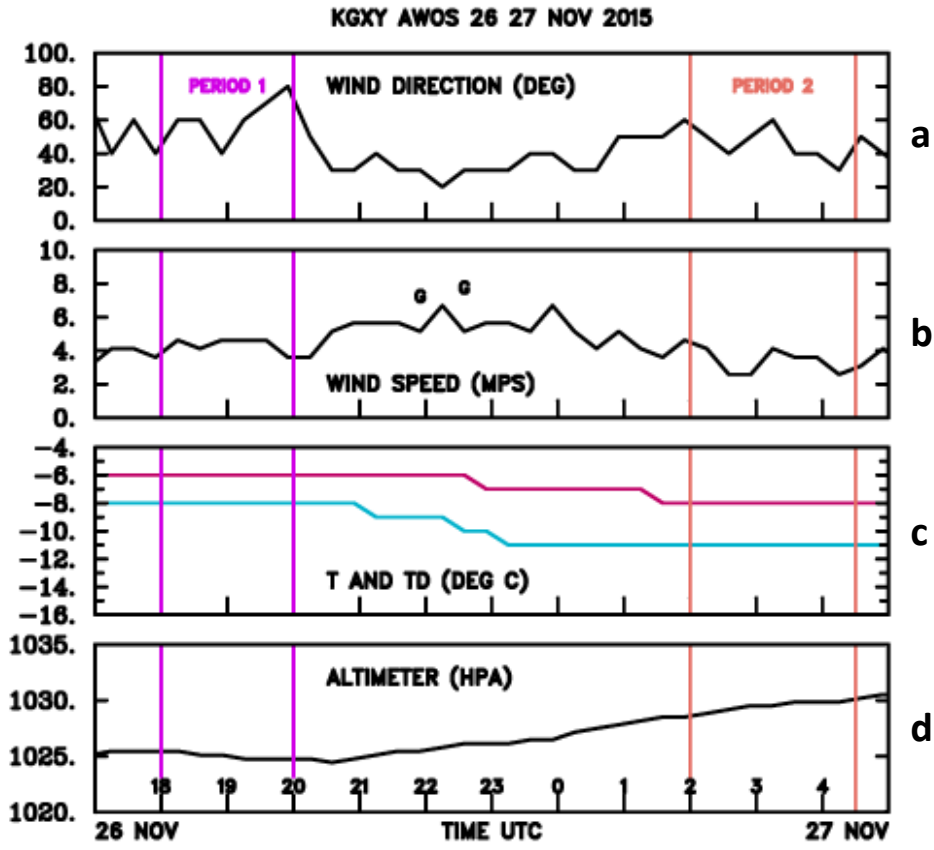
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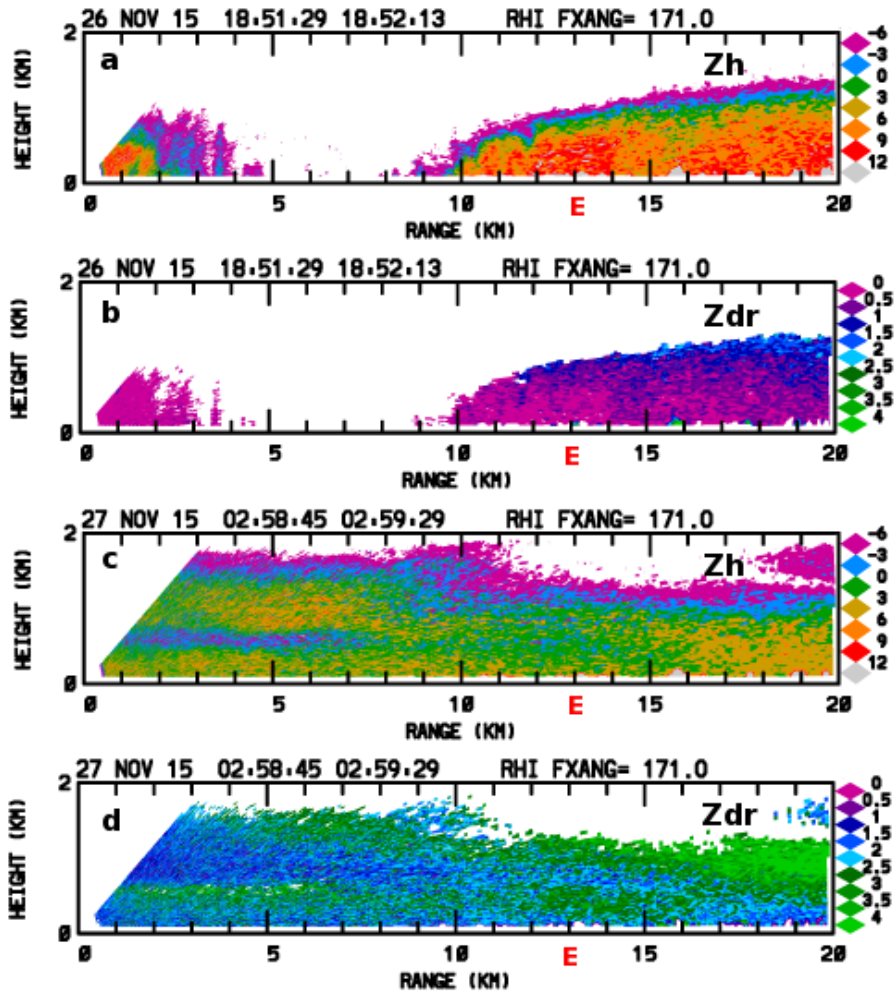


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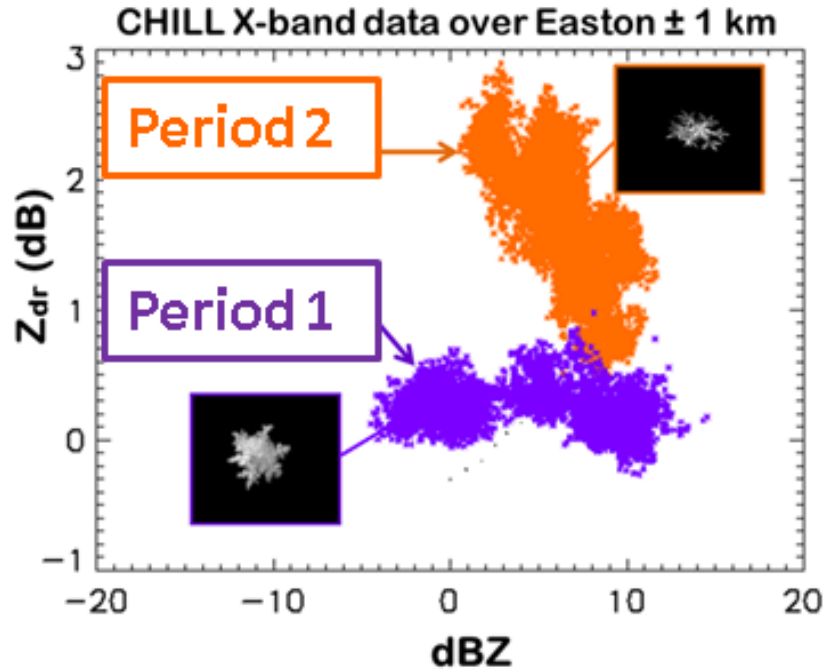
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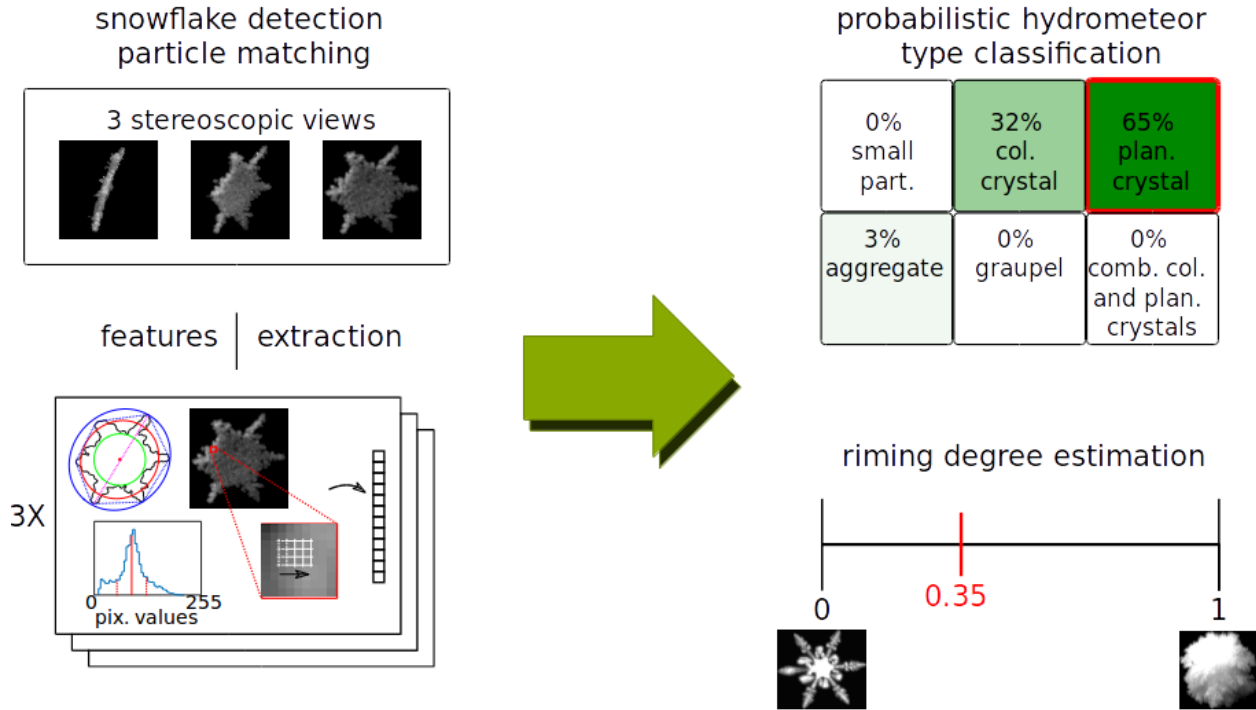
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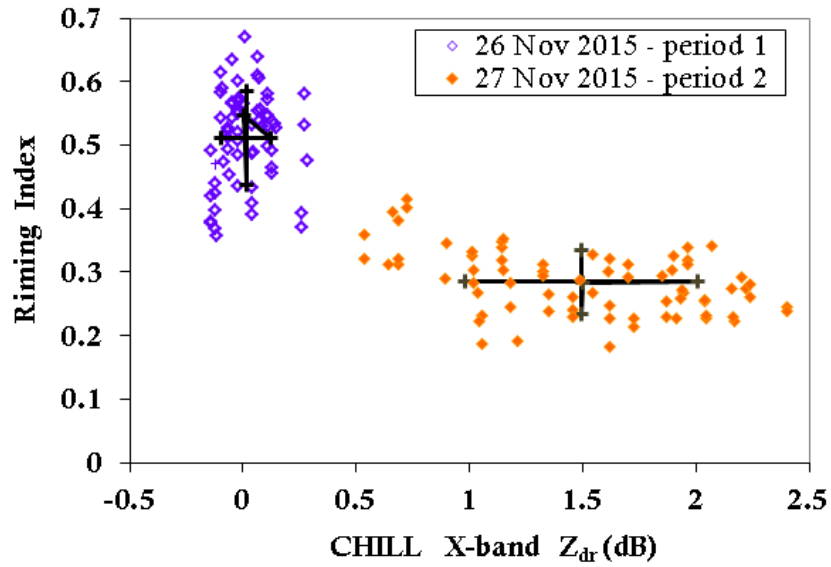
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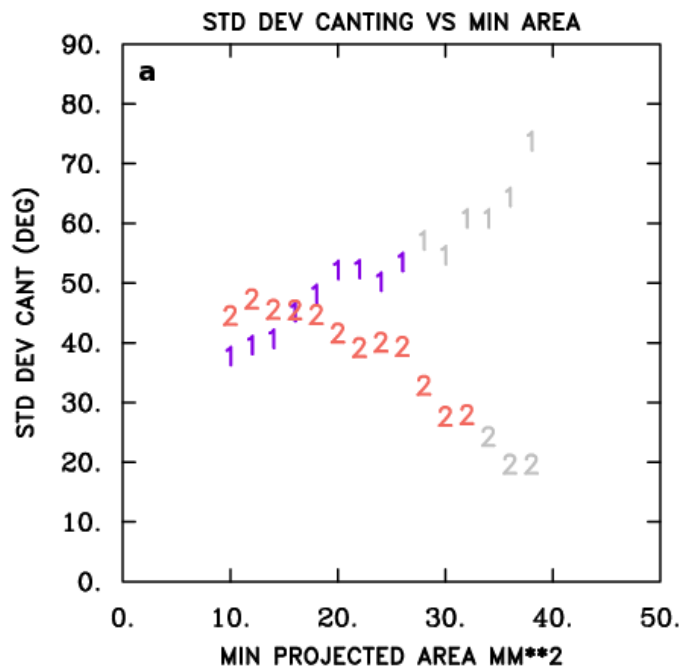
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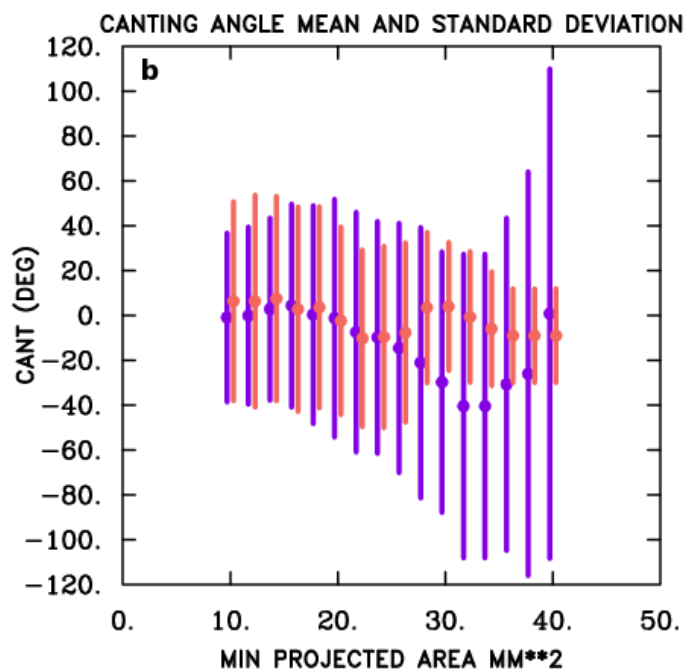
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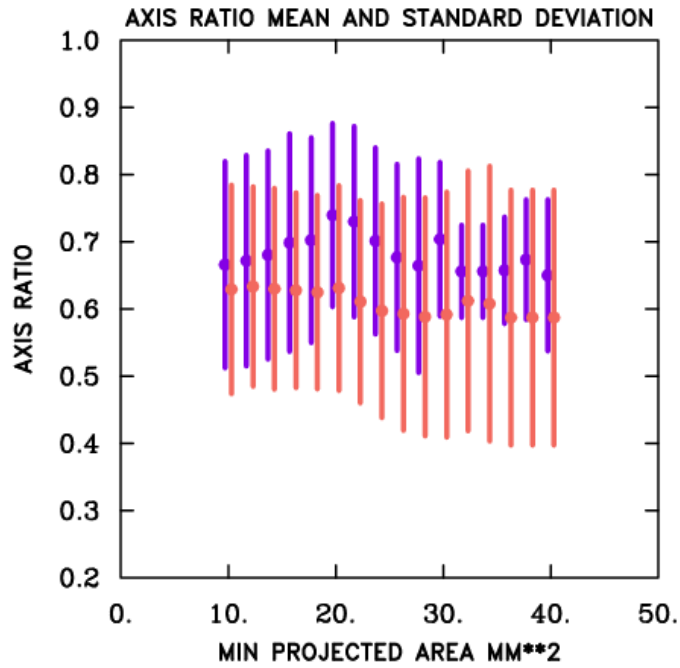


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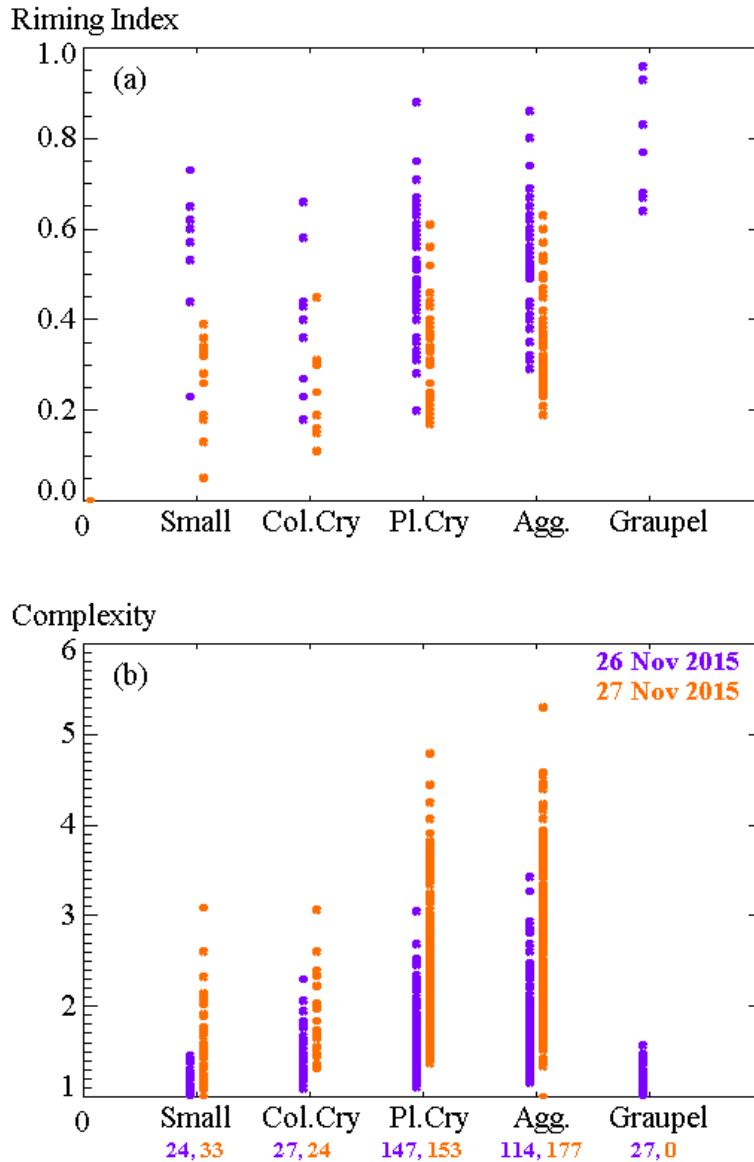
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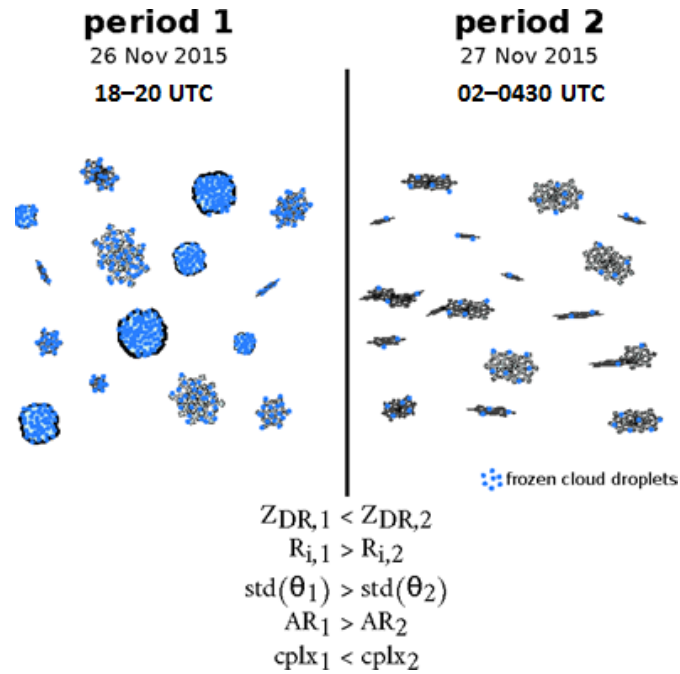
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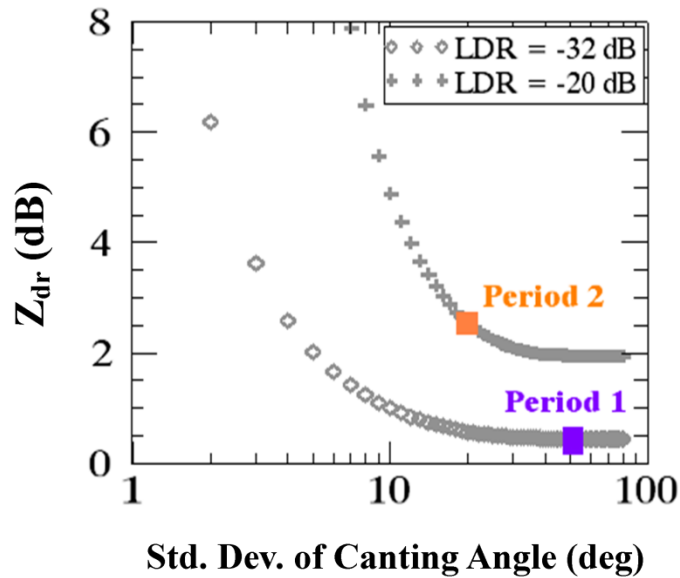
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